

# Multigap superconductivity in $\text{La}_2\text{C}_3$ and $\text{Y}_2\text{C}_3$

J. Akimitsu

*Department of Physics and Mathematics, Aoyama Gakuin University*

The discovery of high-temperature superconductivity in  $\text{MgB}_2$  has activated renewed interest in other boride and carbide superconductors. Among such compounds, sesquicarbides ( $\text{Ln}_2\text{C}_3$ ,  $\text{Ln} = \text{La, Y}$ ) show relatively high critical temperatures and their  $T_c$  values strongly depend on their carbon composition.

Recently, we have found a new superconducting phase in  $\text{Y}_2\text{C}_3$ , exhibiting a much higher  $T_c$  ( $T_c \sim 18$  K). A recent study on the temperature dependence of the nuclear spin lattice relaxation rate in  $\text{Y}_2\text{C}_3$  has suggested the multiple superconducting gaps with s-wave symmetry. While a similar electronic structure would be expected for a  $\text{La}_2\text{C}_3$ , the specific measurement of this material suggests single gap superconductivity.

The  $\mu\text{SR}$  technique is a powerful microscopic tool to determine the existence of two gap state. In this conference, we present result of  $\mu\text{SR}$  measurement on polycrystalline samples of  $\text{La}_2\text{C}_3$  ( $T_c \sim 11$  K) and  $\text{Y}_2\text{C}_3$  ( $T_c \sim 15$  K), where a clear sign of double gap superconductivity was observed in the temperature dependence of muon depolarization rate. The muon depolarization rate ( $\sigma_v(T)$ ) exhibits a characteristic temperature dependence that can be perfectly described by a phenomenological double-gap model for nodeless superconductivity (see Fig. 1). While the magnitude of two gaps is similar between  $\text{La}_2\text{C}_3$  and  $\text{Y}_2\text{C}_3$ , a significant difference in the interband coupling between those two cases was clearly observed in the behavior of  $\sigma_v(T)$ .

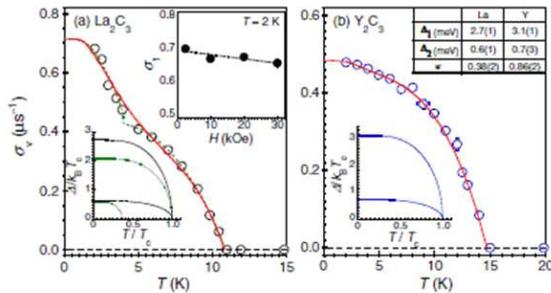


Fig. 1 Temperature dependence of the muon spin relaxation rate for (a)  $\text{La}_2\text{C}_3$  at 2.5 kOe and (b)  $\text{Y}_2\text{C}_3$  at 5.0 kOe. Error bars (not shown) are smaller than the symbol size. Solid and dashed curves indicate the result of the fitting analysis using the double-gap model described in the text. Insets show the relaxation rate in the superconducting state ( $\sigma_1$ ) as a function of the magnetic field for  $\text{La}_2\text{C}_3$  (a) and the order parameters  $[\Delta(T)/k_B T_c]$  for the respective cases.

	$\text{La}_2\text{C}_3$	$\text{Y}_2\text{C}_3$
Transverse field (kOe)	2.5	5.0
$T_c$ (K)	10.9(1)	14.7(2)
$\sigma_v(0)$ ( $\mu\text{s}^{-1}$ )	0.71(3)	0.48(2)
$\lambda(0)$ (Å)	3800(100)	4600(100)
$w$	0.38(2)	0.86(2)
$\Delta_1(0)$ (meV)	2.7(1)	3.1(1)
$\Delta_2(0)$ (meV)	0.6(1)	0.7(3)
$2\Delta_1/k_B T_c$	5.6(3)	4.9(3)
$2\Delta_2/k_B T_c$	1.3(3)	1.1(5)

TABLE I. Superconducting properties of  $\text{La}_2\text{C}_3$  and  $\text{Y}_2\text{C}_3$  determined from the present experiment, where those obtained from the double-gap analysis correspond to the solid curves in Fig. 1.